An Introduction to Parallel-in-time methods

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Paris, 16th of May 2023

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Introduction to Parallel-in-time methods

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What is a Parallel-in-Time method?

Methods that compute solutions at a further time step before that the solution at a closer time step has been computed.

- Usually iterative methods
- Usually comes at the expense of additional work

Why are Parallel-in-Time methods interesting?

- Problems that need a solution by a certain deadline
- Problems that are very long in time

Traditionally, Parallel-in-Time methods are classified in the following categories:

- Shooting type methods
 - Parareal

2 Multigrid methods

- Multigrid Reduction-in-Time
- Space-Time Multigrid
- Oomain Decomposition methods
 - Schwarz Waveform Relaxation
- 4 Direct methods
 - ParaExp

Space-time decomposition in ...



- 50 Years of Time Parallel Time Integration (2015) M. J. Gander
- Applications of time parallelization (2020) B. W. Ong and J. B. Schroder
- Multigrid methods with space-time concurrency (2017)
 R. D. Falgout, S. Friedhoff, Tz. V. Kolev, S. P. MacLachlan, J. B. Schroder, S. Vandewalle

We will solve the 1d heat equation (c = 1)

$$\begin{cases} u_t(x,t) = u_{xx}(x,t) , & (x,t) \in (0,L) \times (0,T) , \\ u(0,t) = u(L,t) = 0 , & t \in (0,T) , \\ u(x,0) = u_0(x) , & x \in (0,L) . \end{cases}$$

We will discretize it using finite differences in space,

$$\left\{ egin{array}{l} v_t(t) = rac{1}{\Delta x^2} \, L \, v(t) \;, \quad t \in (0, \, T) \ v(0) = u_0 \;. \end{array}
ight.$$

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			1	-2/

Given the problem

$$\left\{ egin{aligned} &v_t(t)=rac{1}{\Delta x^2}\,L\,v(t)\,,\quad t\in(0,\,T)\ &v(0)=u_0\in\mathbb{R}^{n_x}\,. \end{aligned}
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We discretize in time using a Runge-Kutta scheme

$$u_{n+1} = \Phi u_n$$
, $u_0 = v(0) \in \mathbb{R}^{n_x}$, $n = 0, 1, \dots, N_t$

Example: If we use Backward Euler $\Phi = (I - \frac{\Delta t}{\Delta x^2}L)^{-1}$.



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Parareal

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Method introduced in a 3-page paper in 1964 by J. Nievergelt.



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Developments of the method

1989 Bellen and Zenaro: solve $u_{n+1} = \Phi u_n$ using (a variant of) Newton's method.

$$\boldsymbol{u}^{k+1} = \varphi(\boldsymbol{u}^k) + \Delta \varphi(\boldsymbol{u}^k)(\boldsymbol{u}^{k+1} - \boldsymbol{u}^k) ,$$

where $\varphi(\boldsymbol{u}) = \begin{pmatrix} u_0 - v(0) \\ u_1 - \Phi u_0 \\ \vdots \\ u_N - \Phi u_{N-1} \end{pmatrix}$ and $\boldsymbol{u} = \begin{pmatrix} u_0 \\ u_1 \\ \vdots \\ u_N \end{pmatrix}$

- Quadratic convergence
- Finite-time convergence

1993 Chartier and Philippe: Noticed that it isn't effective for all kinds of problems.

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Parareal¹

Consider a coarsening factor of m (fine/coarse grid). Let $F = \Phi^m$ and G a cheap approximation of F.

Idea: Approximate $\Delta \varphi$ by the finite difference $(G U_n^{k+1} - G U_n^k)/(U_n^{k+1} - U_n^k)$.

Initialization:

$$\begin{cases} U_0^0 = u_0 \\ U_{n+1}^0 = G \ U_n^0 \ , \quad n = 0, \dots, N_t - 1 \ . \end{cases}$$

Parareal iteration: for $k = 0, \ldots, K - 1$

$$\begin{cases} U_0^{k+1} = u_0 , \\ U_{n+1}^{k+1} = F U_n^k + G U_n^{k+1} - G U_n^k , & n = 0, \dots, N_t - 1 . \end{cases}$$

¹Lions, Maday, Turinici (2001)

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Illustration of Parareal

Initialization:

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Illustration of Parareal

Parareal iteration: for $k = 0, \ldots, K - 1$

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Introduction to Parallel-in-time methods

- Nievergelt's method is too expensive for large problems
- Parallelism in Parareal is achieved though additional work of a coarse operator
- For Parareal to be effective, the coarse operator needs to:
 - be a good approximation of the fine
 - be much cheaper to compute

- Résolution d'EDP par un schéma en temps "pararéel" (2001) J.-L. Lions, Y. Maday, G. Turinici
- A Micro-Macro Parareal Algorithm: Application to Singularly Perturbed Ordinary Differential Equations (2013)
 F. Legoll, T. Lelièvre, G. Samaey
- PARAOPT: A Parareal Algorithm for Optimality Systems (2020) M. J. Gander, F. Kwok, J. Salomon
- Low-rank Parareal: a low-rank parallel-in-time integrator (2023)
 B. Carrel, M. J. Gander, B. Vandereycken

Multigrid Reduction-in-Time

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All-at-once system

We are interested in solving the time stepping scheme

$$u_{n+1} = \Phi u_n$$
, $u_0 = u_0$, $n = 0, 1, \dots, N_t - 1$.

This can be written as the following system to solve



 \rightarrow It can be solved using Multigrid!

The Multigrid Reduction-in-Time algorithm² (MGRIT) follows the same steps as a traditional Multigrid algorithm:

- Pre-smoothing: FCF-relaxation
- Omputation of the residual and restriction: Injection
- Coarse grid solve
- Prolongation and correction: Ideal prolongation
- Sost-smoothing: None

²Falgout, Friedhoff, Kolev, MacLachlan, Schroder (2014)

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C-points points that belong both to the coarse and fine grids F-points the others -relaxation update C-points

F-relaxation update F-points

Link to Parareal: $F = \Phi^m$

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Smoothing: What is FCF-relaxation?



C-points points that belong both to the coarse and fine grids F-points the others C-relaxation update C-points F-relaxation update F-points

Link to Parareal: $F = \Phi^m$

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Smoothing: What is FCF-relaxation?



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Link to Parareal: $F = \Phi^m$

Transfer operators in MGRIT



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Parareal and MGRIT are related

Recall (Parareal)

$$\begin{cases} U_0^{k+1} = u_0 & k = 0, \dots, K, \\ U_{n+1}^{k+1} = F U_n^k + G U_n^{k+1} - G U_n^k & n = 0, \dots, N_t - 1, \quad k = 0, \dots, K. \end{cases}$$

Theorem (Gander, Kwok, Zhang, 2018)

The two-level MGRIT algorithm with FCF-relaxation computes the same iterations as the Parareal algorithm using generous overlap of one coarse time interval

Parareal (MGRIT with F-relaxation)



MGRIT with FCF-relaxation



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How do Parareal and MGRIT converge?

The convergence of Parareal: from the iteration

$$U_{n+1}^{k+1} = F U_n^k + G U_n^{k+1} - G U_n^k, \quad n = 1, \dots, n_t ,$$

the error, $\epsilon_n^k := u_n - U_n^k$, can be computed as

$$\epsilon_{n+1}^{k+1} = F \, \epsilon_n^k + G \, \epsilon_n^{k+1} - G \, \epsilon_n^k \; .$$

In turn, it can be bounded as

$$\|\epsilon_{n+1}^{k+1}\| =: e_{n+1}^{k+1} \leq \underbrace{\|\mathcal{F} - G\|}_{=\alpha} e_n^k + \underbrace{\|G\|}_{=\beta} e_n^{k+1}.$$

We thus only need to solve the iteration

$$e_{n+1}^{k+1} = \alpha \, e_n^k + \beta \, e_n^{k+1}$$

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How do Parareal and MGRIT converge?

The convergence of Parareal: from the iteration

$$U_{n+1}^{k+1} = F U_n^k + G U_n^{k+1} - G U_n^k, \quad n = 1, \dots, n_t ,$$

the error, $\epsilon_n^k := u_n - U_n^k$, can be computed as

$$\epsilon_{n+1}^{k+1} = F \, \epsilon_n^k + G \, \epsilon_n^{k+1} - G \, \epsilon_n^k \, .$$

In turn, it can be bounded as

$$\|\epsilon_{n+1}^{k+1}\| =: e_{n+1}^{k+1} \leq \underbrace{\|\mathcal{F} - G\|}_{=lpha} e_n^k + \underbrace{\|G\|}_{=eta} e_n^{k+1}.$$

We thus only need to solve the iteration

$$\mathbf{e}_{n+1}^{k+1} = \alpha \, \mathbf{e}_n^k + \beta \, \mathbf{e}_n^{k+1}$$

The iteration $e_{n+1}^k = \alpha e_n^{k-1} + \beta e_n^k$ can be written in matrix form as

$$\begin{pmatrix} I & & & \\ -\beta & I & & \\ & \ddots & \ddots & \\ & & -\beta & I \end{pmatrix} \begin{pmatrix} e_0^k \\ e_1^k \\ \vdots \\ e_{n_t}^k \end{pmatrix} = \begin{pmatrix} 0 & & & \\ \alpha & 0 & & \\ & \ddots & \ddots & \\ & & \alpha & 0 \end{pmatrix} \begin{pmatrix} e_0^{k-1} \\ e_1^{k-1} \\ \vdots \\ e_{n_t}^{k-1} \end{pmatrix}$$

Lemma (Recurrence solving)

Assuming that α and β are scalars, the error at step k is given by

$$\boldsymbol{e}^{k} = \boldsymbol{M}(\beta) \left(\boldsymbol{I}_{n_{t}} \otimes \alpha
ight) \boldsymbol{e}^{k-1} = \ldots = \boldsymbol{M}(\beta)^{k} \left(\boldsymbol{I}_{n_{t}} \otimes \alpha^{k}
ight) \boldsymbol{e}^{0} ,$$

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How do Parareal and MGRIT converge?

We then want to bound

$$oldsymbol{e}^k = oldsymbol{M}(eta)^k \left(oldsymbol{I}_{n_t} \otimes lpha^k
ight) oldsymbol{e}^0 \; .$$

Linear Bound [Lemma 4.4, Gander, Vandewalle, 2007]

$$\|M(eta)^k\|_\infty \leq \|M(eta)\|_\infty^k = \left(rac{1-|eta|^{n_t}}{1-|eta|}
ight)^k$$

Superlinear Bound [Lemma 4.3, Gander, Vandewalle, 2007]

$$\|M(\beta)^{k}\|_{\infty} = \sum_{i=0}^{n_{t}-k} \binom{i+k-1}{k-1} |\beta|^{i} = \frac{1}{(k-1)!} \sum_{i=0}^{n_{t}-k} \left[\prod_{l=1}^{k-1} (i+l)\right] |\beta|^{i}$$

M. J. Gander, S. Vandewalle. "Analysis of the Parareal Time-Parallel Time-Integration Method", 2007.

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Linear Bound [Lemma 4.4, Gander, Vandewalle, 2007]

$$\|M(eta)^k\|_\infty \leq \|M(eta)\|_\infty^k = \left(rac{1-|eta|^{n_t}}{1-|eta|}
ight)^k$$

Explicit Superlinear Bound [Lemma 4.4, Gander, Vandewalle, 2007]

If $|\beta| < 1$, then

$$\|M(\beta)^k\|_{\infty} \leq \binom{n_t}{k} = \frac{1}{k!} \prod_{l=0}^{k-1} (n_t - l).$$

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Bound on the eigenvalues³

Assume λ_F and λ_G are the eigenvalues of F and G. Let $\alpha = \lambda_F - \lambda_G$ and $\beta = \lambda_G$, then

$$\|M(\beta)(I_{n_t} \otimes \alpha)\|_{\infty} \leq |\lambda_F - \lambda_G| \frac{1 - |\lambda_G|^{n_t}}{1 - |\lambda_G|}$$



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Bound on the eigenvalues³

Assume λ_F and λ_G are the eigenvalues of F and G. Let $\alpha = \lambda_F - \lambda_G$ and $\beta = \lambda_G$, then

$$\|M(eta)(I_{n_t}\otimeslpha)\|_\infty\leq |\lambda_{\sf F}-\lambda_{\sf G}|\,rac{1-|\lambda_{\sf G}|^{n_t}}{1-|\lambda_{\sf G}|}$$



³Dobrev, Kolev, Petersson, Schroder (2017)

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Convergence of Parareal and MGRIT

Initial guess: $u_0(x) = \sin(\pi x)$, $x \in [0, 1]$. $n_x = 100$, $n_t = 256$, T = 0.02 for heat and T = 1 for advection. SDIRK2 in time and second order centered for heat and upwind for advection.



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Convergence of Parareal and MGRIT



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- MGRIT is a multilevel version of Parareal
- MGRIT usually uses FCF-relaxation, whereas Parareal uses F-relaxation
- From superlinear bound: superlinear + finite time convergence
- From linear bound: eigenvalues that are close to unity need to be well approximated.

- Analysis of the Parareal Time-Parallel Time-Integration Method (2007) M. J. Gander, S. Vandewalle
- Parallel Time Integration with Multigrid (2014) R. D. Falgout, S. Friedhoff, Tz. V. Kolev, S. P. MacLachlan, J. B. Schroder
- Two-Level Convergence Theory for Multigrid Reduction in Time (MGRIT) (2017) V. A. Dobrev, Tz. V. Kolev, N. A. Petersson, J. B. Schroder
- Scheduling of tasks in the Parareal algorithm (2011) E Aubanel
- A Unified Analysis Framework for Iterative Parallel-in-Time Algorithms (2022) M. J. Gander, T. Lunet, D. Ruprecht, R. Speck

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Space-Time Multigrid

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The time-stepping procedure

$$u_{n+1} = \Phi u_n$$
, $u_0 = u_0$, $n = 0, 1, \dots, N_t - 1$,

was formulated as a all-at-once system for MGRIT. For the Space-Time algorithm, we consider a different formulation.

Separate the operator $\Phi = Q^{-1}P$ where Q and P are the implicit and explicit parts of Φ . In turn, it can be written as the all-at-once system,



Idea: Apply Multigrid on the space-time problem.

- Pre-smoothing: ν_1 damped Jacobi iterations
- Output Computation of the residual and restriction: Full-weighting
- Coarse grid solve
- Prolongation and correction: Linear interpolation
- Post-smoothing: ν_2 damped Jacobi iterations

Reminder: The point-wise recall the damped Jacobi smoother is given by

$$oldsymbol{v}^{k+1} = oldsymbol{v}^k + \omega D^{-1} [oldsymbol{b} - Aoldsymbol{v}^k]$$

where D is the diagonal matrix such that $D_{ii} = A_{ii}$.

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⁴Horton, Vandewalle (1995)

Transfer operators in STMG



$$\longrightarrow$$
 $\times 1$ \longrightarrow $\times 0.5$ \longrightarrow $\times 0.25$

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Space-Time Multigrid with point-wise smoother



Figure: Convergence factor of STMG with point-wise smoother from Horton, Vandewalle (1995)

Goal: Be able to always smooth in the time dimension.

Consider the damped block-Jacobi smoother

$$\boldsymbol{v}^{k+1} = \boldsymbol{v}^k + \omega D^{-1} [\boldsymbol{b} - A \boldsymbol{v}^k] ,$$

with D a block diagonal matrix with blocks Q and $\omega \in [0,2]$ a damping parameter.

 \rightarrow Always allows for coarsening in time

⁵Gander, Neumüller (2016)

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STMG has good scaling properties

Weak and Strong scaling results: 3D Heat Equation

	Weak scaling						Strong scaling				
cores	n _t	dof	iter	time	fwd. subs.	n _t	dof	iter	time		
1	2	59,768	7	28.8	19.0	512	15,300,608	7	7,635.2		
2	4	119,536	7	29.8	37.9	512	15,300,608	7	3,821.7		
4	8	239,072	7	29.8	75.9	512	15,300,608	7	1,909.9		
8	16	478,144	7	29.9	152.2	512	15,300,608	7	954.2		
16	32	956,288	7	29.9	305.4	512	15,300,608	7	477.2		
32	64	1,912,576	7	29.9	613.6	512	15,300,608	7	238.9		
64	128	3,825,152	7	29.9	1,220.7	512	15,300,608	7	119.5		
128	256	7,650,304	7	29.9	2,448.4	512	15,300,608	7	59.7		
256	512	15,300,608	7	30.0	4,882.4	512	15,300,608	7	30.0		

Table: Vulcan BlueGene /Q Supercomputer in Livermore (Martin Neumüller)

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STMG has good scaling properties

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2	4	119,536	7	29.8	37.9	512	15,300,608	7	3,821.7		
4	8	239,072	7	29.8	75.9	512	15,300,608	7	1,909.9		
8	16	478,144	7	29.9	152.2	512	15,300,608	7	954.2		
16	32	956,288	7	29.9	305.4	512	15,300,608	7	477.2		
32	64	1,912,576	7	29.9	613.6	512	15,300,608	7	238.9		
64	128	3,825,152	7	29.9	1,220.7	512	15,300,608	7	119.5		
128	256	7,650,304	7	29.9	2,448.4	512	15,300,608	7	59.7		
256	512	15,300,608	7	30.0	4,882.4	512	15,300,608	7	30.0		

Table: Vulcan BlueGene /Q Supercomputer in Livermore (Martin Neumüller)

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The convergence of STMG depends on the ratio $\sigma = \Delta t / \Delta x^2$

Let μ be the smoothing factor associated to the block-Jacobi smoother obtained by Local Fourier Analysis (LFA),



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- STMG has very good scaling properties
- The convergence of STMG depends on the ratio $\Delta t/\Delta x^2$
- STMG with block-smoothing allows to always coarsen in time

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- A Space-Time Multigrid for Parabolic Partial Differential Equations (1995) G. Horton, S. Vandewalle
- Analysis of a New Space-Time Parallel Multigrid Algorithm for Parabolic Problems (2016) M. J. Gander, M. Neumüller
- An Optimized Space-Time Multigrid Algorithm for Parabolic PDEs (2023, in review) B. Chaudet-Dumas, M. J. Gander, A. P.

Schwarz Waveform Relaxation

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Schwarz Waveform Relaxation

We will solve the 1D heat equation

$$\left\{egin{aligned} & u_t(x,t) = u_{xx}(x,t) \;, & (x,t) \in \Omega imes (0,T) \;, \ & u(x,t) = 0 \;, & x \in \partial \Omega, \; t \in (0,T) \;, \ & u(x,0) = u_0(x) \;, & x \in \Omega \;. \end{aligned}
ight.$$

Partition the domain $\Omega=\Omega_1\cup\Omega_2.$ Define for $i\in\{0,1\},$

$$\begin{cases} \partial_t u_i^{k+1}(x,t) = \partial_{xx} u_i^{k+1}(x,t) & (x,t) \in \Omega_1 \times (0,T] \\ u_i^{k+1}(x,0) = u_0(x) & x \in \Omega_i \\ u_i^{k+1}(x,t) = 0 & x \in \partial\Omega, \ t \in (0,T] \\ u_i^{k+1}(\Gamma_i,t) = u_{1-i}^k(\Gamma_i,t) & t \in (0,T] \end{cases}$$

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Schwarz Waveform Relaxation

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$$\left\{ egin{array}{ll} u_t(x,t) = u_{xx}(x,t) \;, & (x,t) \in \Omega imes (0,T) \;, \ u(x,t) = 0 \;, & x \in \partial \Omega, \; t \in (0,T) \;, \ u(x,0) = u_0(x) \;, & x \in \Omega \;. \end{array}
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0 Γ_1 Ω_2 u_0 Ω_1 Γ_2 Ω Paris, 16th of May 2023 39 / 54

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Goal: Compute the solution for the 1D heat equation with Dirichlet boundary conditions with $u_0 = 20$ and initial guess $u_0 = 0$.

On the right, first 6 iterations of SWR at time T = 0.5.



Convergence is superlinear⁶



Optimized Schwarz Waveform Relaxation

Recall: Schwarz wave relaxation is given for $i \in \{0, 1\}$ by

$$\begin{cases} \partial_t u_i^{k+1}(x,t) = \partial_{xx} u_i^{k+1}(x,t) & (x,t) \in \Omega_1 \times (0,T] \\ u_i^{k+1}(x,0) = u_0(x) & x \in \Omega_i \\ u_i^{k+1}(x,t) = 0 & x \in \partial\Omega, \ t \in (0,T] \\ u_i^{k+1}(\Gamma_i,t) = u_{1-i}^k(\Gamma_i,t) & t \in (0,T] \end{cases}$$

Goal: Do SWR without overlap.

$$\begin{cases} \partial_t u_i^{k+1}(x,t) = \partial_{xx} u_i^{k+1}(x,t) & (x,t) \in \Omega_1 \times (0,T] \\ u_i^{k+1}(x,0) = u_0(x) & x \in \Omega_i \\ u_i^{k+1}(x,t) = 0 & x \in \partial\Omega, \ t \in (0,T] \\ (\partial_x + p) u_i^{k+1}(\Gamma_i,t) = (\partial_x + p) u_{1-i}^k(\Gamma_i,t) & t \in (0,T] \end{cases}$$

Solution: Change transmission conditions to Robin type and optimize for p_{\odot} ,

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Optimizing Optimized Schwarz Waveform Relaxation



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- Schwarz Wave Relaxation allows parallelism in time by cutting the spatial domain in two (or more) subdomains
- Schwarz Wave Relaxation has superlinear convergence
- Optimized Schwarz Waveform Relaxation allows to have no overlap between domains.

- The waveform relaxation method for time-domain analysis of large scale integrated circuits (1983)
 - E. Lelarasmee, A. E. Ruehli, and A. L. Sangiovanni-Vincentelli.
- Overlapping Schwarz of linear and non-linear Parabolic problems (1996) M. J. Gander
- Dirichlet-Neumann and Neumann-Neumann waveform relaxation algorithms for parabolic problems (2013)
 M. J. Gander, F. Kwok, B. C. Mandal
- Work by M. Gander, L. Halpern, V. Martin, F. Nataf.



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Introduced by Gander and Güttel in 2013. Let $A \in \mathbb{R}^{n \times n}$ and f a non-linear function.

$$u'(t) = A u(t) + f(t), \quad t \in (0, T), \quad u(0) = u_0.$$

Observation

A homogenous problem can be integrated much faster than an inhomogenous problem.

Separate u = v + w, where

$$\mathbf{v}'(t) = A \mathbf{v}(t), \quad \mathbf{v}(0) = u_0,$$

 $\mathbf{w}'(t) = A \mathbf{w}(t) + f(t), \quad \mathbf{w}(0) = 0.$

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ParaExp — illustration



Figure: Illustration of ParaExp from the original article (Gander, Güttel, 2013)

- On [T_n, T_{n+1}]: Solve with serial integrator
 On [T_n, T]: Solve with near optimal exponential integrator
- **③** Get the final solution by superposition of solutions

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Introduction to Parallel-in-time methods

Question: If A is very big, how to compute efficiently $\exp(tA)$?

Idea: Compute exp(tA)v directly for some initial condition v using Krylov methods.

Generate the Krylov space \$\mathcal{K}_m = span{\$v, Av, \ldots, A^{m-1}v\$}\$.
 Use the Arnoldi method.

- **2** Get an orthonormal basis V_m of \mathcal{K}_m .
 - Use a modified Gram-Schmidt process.

Compute H_m the projection of A on the space K_m with respect to the basis V_m.
 The matrix H_m is upper Hessenberg and verifies H_m = V_m^T AV_m.

• Get the approximation $e^{tA} \approx \beta V_m e^{tH_m} e_1$ (good even for small m)

• Apply traditional techniques to compute $e^{tH_m}e_1$.

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Question: If A is very big, how to compute efficiently exp(tA)?

Idea: Compute exp(tA)v directly for some initial condition v using Krylov methods.

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- **2** Get an orthonormal basis V_m of \mathcal{K}_m .
 - Use a modified Gram-Schmidt process.
- **③** Compute H_m the projection of A on the space \mathcal{K}_m with respect to the basis V_m .
 - The matrix H_m is upper Hessenberg and verifies $H_m = V_m^\top A V_m$.
- Get the approximation $e^{tA} \approx \beta V_m e^{tH_m} e_1$ (good even for small m)
 - Apply traditional techniques to compute $e^{tH_m}e_1$.

Consider the non-linear problem, with $B : \mathbb{R}^n \to \mathbb{R}^n$ a non-linear application,

$$u'(t) = A u(t) + B(u(t)) + f(t), t \in [0, T],$$
 $u(0) = u_0.$

Problem: Applying the splitting u = v + w does not verify the original equation.

$$v'(t) = A v(t),$$

 $v'(t) = A w(t) + B(w(t)) + f(t),$
 $v(0) = u_0,$
 $w(0) = 0.$

Non-linear ParaExp

Consider the following iterative method: for k = 0 initialize

$$u_n^0 = w_n^0 = 0, \qquad n = 0, \dots, N_t - 1.$$

For $k = 1, 2, \ldots$ solve the homogenous problems

$$(w_n^k)'(t) = A w_n^k(t) \qquad t \in [T_{n-1}, T_n]$$
$$w_1^k(T_0) = u_0, \ w_n^k(0) = u_{n-1}^{k-1}(0) - \sum_{j=1}^{n-1} w_j^{k-1}(0)$$

and then solve the non-homogenous problem

$$(u_n^k)'(t) = A u_n^k(t) + B(u_n^k(t)) + f(t),$$
 $t \in [T_{n-1}, T_n]$
 $u_n^k(0) = \sum_{j=1}^n w_j^k(T_{n-1})$

Theorem [Theorem 2, Gander, Güttel, Petcu, 2018]

Let the coarse propagator $G(U_n)$ solve the linear problem

 $u'(t) = A u(t), \quad t \in (T_n, T_{n+1}), \quad u(0) = U_n$

and the fine propagator $F(U_n)$ solve the non-linear problem

$$u'(t) = A u(t) + B(u(t)) + f(t), \quad t \in (T_n, T_{n+1}), \quad u(0) = U_n$$

Then the Parareal iteration computed with those operators is equivalent to the non-linear ParaExp algorithm.
- Direct method for linear problems
- Acceleration of computation of the exponential using Krylov methods
- Iterative method for non-linear problems
- Equivalence of non-linear ParaExp with Parareal

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- ParaExp: a Parallel Integrator for Linear Initial-Value Problems (2013)
 M. J. Gander, S. Güttel
- A Nonlinear ParaExp Algorithm (2018) M. J. Gander, S. Güttel, M. Petcu
- Analysis of Some Krylov Subspace Approximations to the Matrix Exponential (1992)
 Y. Saad